## SOME PROBLEMS TO SOLVE

**A**. By inspection of Maxwell's equations show that the dimension of c is that of length/time. Starting from the equations in SI units, rescale **E** and **B** with dimensionful constants so they get the same dimension.

**B**. Assuming that Stokes' theorem holds in three dimensions prove that Gauss' theorem holds in two dimensions. (Hint: In two dimension there is a unique direction normal to any given vector.)

C. Prove entry 5 in Jackson's list of properties of the delta function (p. 26).

**D**. An analytic function is a complex valued function on a two dimensional plane such that it depends only on the particular combination z = x + iy,  $\Phi(x, y) = \Phi(x + iy) = \Phi(z)$ . Prove that every such function is a solution of the two dimensional Laplace equation.

**E**. Use the chain rule for derivatives to express the two dimensional Laplace operator in polar coordinates  $(x = r \cos \phi, y = r \sin \phi)$ .

**F**. The Poisson equation  $\nabla^2 \Phi = 4\pi G \rho_{\rm m}$ , where  $\rho_{\rm m}$  is the mass density, occurs in Newton's theory of gravity. Show that the potential

$$\Phi(x, y, z) = -\frac{GM}{\sqrt{x^2 + y^2 + (|z| + a)^2}}$$
(1)

comes from an infinitely thin disk of matter (that is, roughly, a spiral galaxy).

**G**. Create an electric field inside an spherical cavity. You can choose the boundary conditions freely. Can you choose them so that a test charge has a stable equilibrium position at the centre of the cavity? Anywhere in the interior of the cavity?

**Ga.** In his eq. (1.58) Jackson includes an integral. He says it is easy to show that it evaluates to  $4\pi$ . Verify this.

**H**. A small charge q sits outside a spherical conductor of large total charge Q and radius R. The charges have the same sign. At (roughly) what distance d from the centre will the charge feel no force?

**Ha**. Two conducting planes are placed at x = 0 and at y = 0. The potential on the planes is kept at zero. Place a single charge somewhere in the positive quadrant of the (x, y)-plane. Use the method of images to calculate the electric potential in that quadrant.

**Hb.** Consider two conducting spheres, radii a and b, distance between centres c, the first sphere at  $\Phi = 0$ , the other at  $\Phi = V$ . Place imaginary charges inside the spheres until the boundary conditions hold, in a stepwise procedure so that the first imaginary charge ensures that  $\Phi = V$  on the second sphere, the second charge is inside the first sphere and compensates for the first charge so that  $\Phi = 0$  on the second sphere. Go on in this way. It is enough to do the first few steps, but after each step you should use a computer to plot the potential at the surface of the first sphere, for suitable values of the parameters. This will give you a feeling for how quickly the procedure converges.

Hc. Derive Jackson's eq. (2.13) from his eq. (2.12), in full detail.

**I**. Compute the first five Legendre polynomials by applying Gram-Schmidt orthogonalization to the polynomials 1, x,  $x^2$ ,  $x^3$ ,  $x^4$ . Also show that you recover the first five Legendre polynomials by expanding the generating function  $g(t,x) = (1 - 2xt + t^2)^{-1/2}$  to fourth order in t.

**J**. Show that the Legendre polynomial  $P_l(x)$  has exactly *l* zeroes in the interval.

**K**. Find the electric field outside a spherical conductor placed in a constant electric field in two ways: using the method of images, and using an expansion in spherical harmonics.

L. In Problem 2.16 Jackson gives an explicit solution for the electric potential inside a square, with a constant charge density there. Is this formula practically useful for plotting the potential? If it is, plot it!

**M**. Compute the first non-vanishing multipole moments for i) two charges q at  $(\pm a, 0, 0)$ , charge -2q at (0, 0, b) ii) four charges q at  $(\pm a, \pm a, 0)$ , two charges -2q at  $(0, 0, \pm b)$ . Check your results.

**N** You have a supply of point charges  $\pm q$ . Place such point charges at the corners of a regular octahedron. You can make the monopole moment vanish. Can you make the dipole moment vanish? The quadrupole moment? Repeat the same exercise for a regular cube.

**O**. For a dipole field, locate those points in space where the field points in a direction orthogonal to the dipole vector.

**P**. A rectangular prism has a square base of area A and a very small height h. Make a cavity of this form inside a dielectric. What is the electric field inside the cavity if the constant **D** in the dielectric is parallel to the height of the cavity? If **D** is orthogonal to the height?

**Q**. Model an atom as a homogeneous sphere of negative charge, and a pointlike nucleus. Assuming a reasonable size, what electric field strength would be needed to shift the nucleus out to one tenth of the radius?

**R**. For pure nitrogen the following data were obtained (in 1934):

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Temperature ( $^{\circ}C$ ):	23.8	23.8	23.8	23.8
Pressure (atm)	1.02	57.5	221.6	1011.6
Density $(kg/m^3)$	1.18	66.04	236.1	578.0
Dielectric constant	1.00052	1.03109	1.11413	1.29633

Check how well this agrees with the Clausius-Mossotti equation.

**S**. Define a vector potential on a region of space strictly outside the z-axis, such that  $\mathbf{A}(\mathbf{x})$  is independent of z, gives a vanishing magnetic field, and cannot be gauge transformed to zero. What is its physical interpretation?

**T**. Solve the Lorentz equation for a charged particle in a homogeneous magnetic field. Based on this solution, would you expect a plasma to be diamagnetic? Paramagnetic?

**U**. The energy of a magnetic dipole placed in a static magnetic field is  $U = -\mathbf{m} \cdot \mathbf{B}$ . There are three cases: permanent magnets ( $\mathbf{m}$  constant), paramagnets and diamagnets ( $\mathbf{m} = k\mathbf{B}$  with k > 0 for paramagnets and k < 0 for diamagnets). For each case investigate whether you can choose  $\mathbf{B}(\mathbf{x})$  such that a stable equilibrium position for the dipole exists.

**V**. Consider the discontinuous function  $\Psi(\mathbf{x}, t) = \Theta(f(\mathbf{x}, t))\psi(\mathbf{x}, t)$ , where  $\Theta$  is the step function (equals 1 if f > 0, 0 if f < 0). Find the conditions on f and  $\psi$  that ensure that we have a solution of the wave equation.

W. Using the retarded Green function solve the inhomogeneous wave equation

$$\Box \psi = \Theta(t) F(t) \delta^{(3)}(x) .$$
<sup>(2)</sup>

("Alice sits at the origin and starts talking at t = 0.") Does the solution make sense?

**X**. In a *d*-dimensional space the Green function for the Laplace equation is (up to a constant)  $G(r) = 1/r^{d-2}$ . Prove that  $\delta(r-t)/r^{d-2}$  is a Green function for the wave equation if and only if d = 3.

**Y**. Draw a map of spacetime in which the *x*-axis and the *t*-axis appear orthogonal. Also draw a lightcone from the origin. Now perform a Lorentz boost.

Draw the x'-axis and the t'-axis on the map. How do they relate to the light cone?

**Z**. Consider a scalar field having the form

$$\phi(x) = \frac{1}{r^2}$$
,  $r^2 = x^2 + y^2 + z^2$ . (3)

Perform a Lorentz boost in the t - x-plane, and express the new function  $\phi'$  that you obtain in this way as a function of the coordinates (t, x, y, z). What does the new function look like?

Å. Consider the electromagnetic field from a point charge at rest at the origin,

$$E_i(x) = \frac{1}{r^3} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \qquad B_i(x) = 0 .$$
(4)

Perform a Lorentz boost in the t - x-plane. Compute the electromagnetic field you obtain, and express it as a function of the coordinates (t, x, y, z).

Ä. Repeat exercise Å, but now for the electromagnetic field

$$E_2 = \cos(t - x)$$
  $B_3 = \cos(t - x)$ , (5)

all other components vanishing.

**Äa.** Prove that Maxwell's equations are not invariant under the Galilei transformation  $t \to t' = t$ ,  $\mathbf{x} \to \mathbf{x}' = \mathbf{x} + \mathbf{v}t$ .

 $\ddot{\mathbf{O}}$ . Use the expression for the energy density of the electromagnetic field to argue that the electric field in a wave sent out in all directions by a localized source falls off like 1/r.